



Thermomagnetic-instability-induced mechanical response in fully magnetized high-temperature superconducting bulk during pulse-field magnetization.

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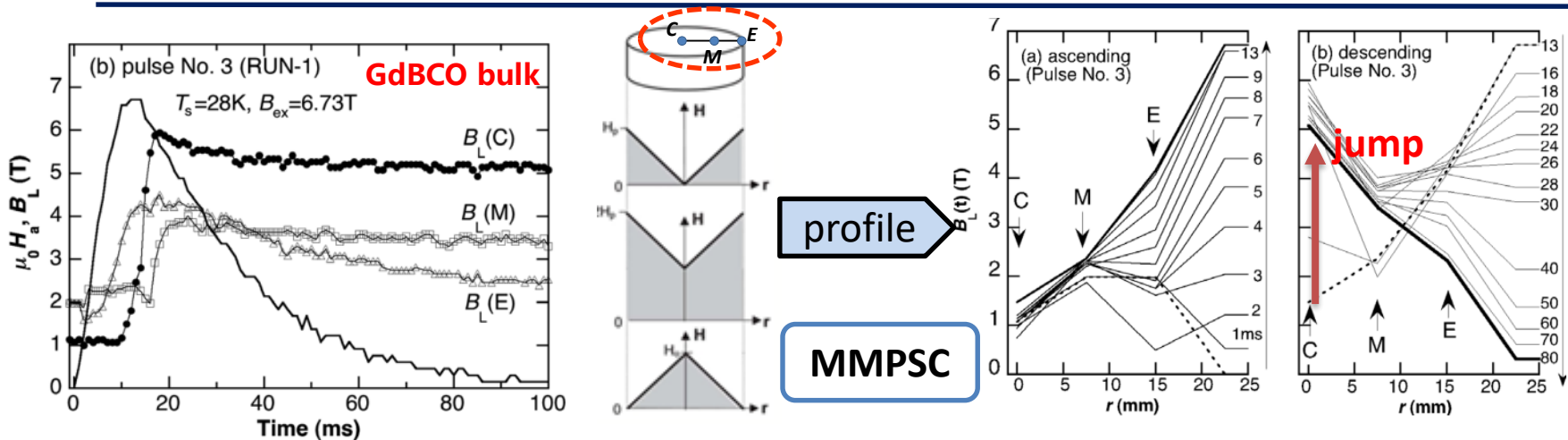
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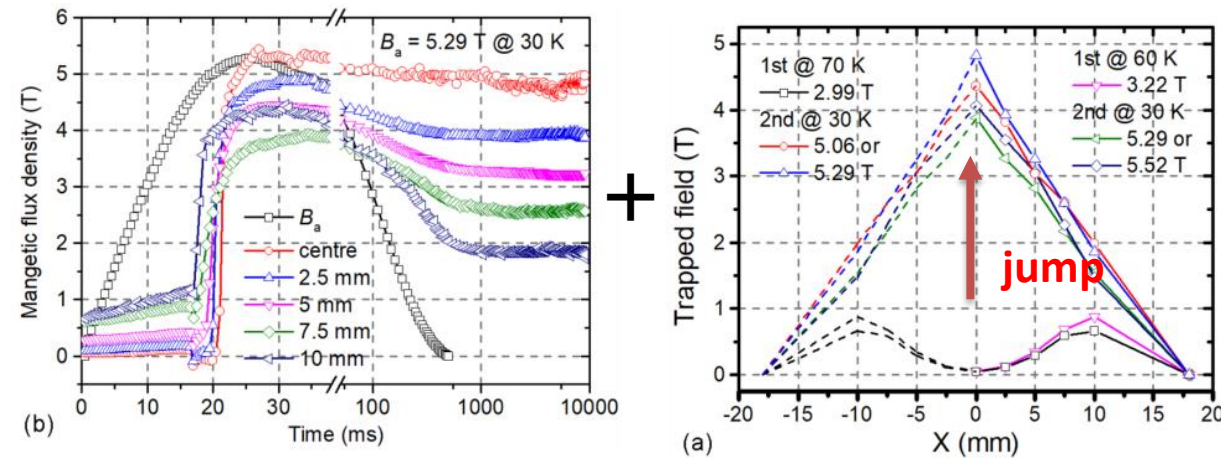
introduction

Flux jump in PFM technique



The trapped field of >5 T achieved at 30 K

Magnetic distribution during flux jump



Flux jump is an essential condition for full magnetization

Variation of magnetic flux density on sample surface

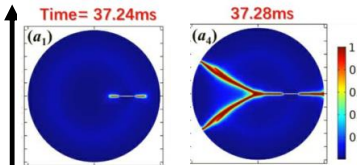
Magnetic distribution during flux jump

Two-step PFM method

Research on the mechanical stability during PFM

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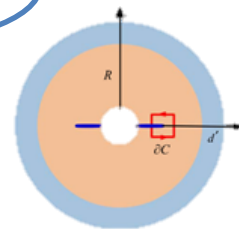
GdBCO bulk



crack
propagati
on and
nucleatio
n

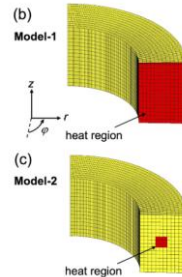
Effect of
crack, void
problem

Low stress
distribution



years

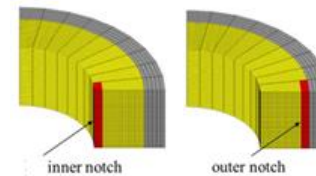
Ring bulk



thermal stress
induced by
local heat
generation

Effect of
Notch to
flux jump

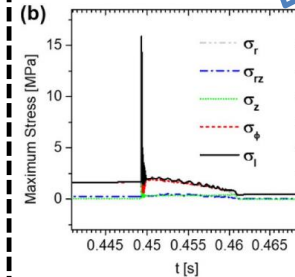
mechanical
fracture in
inner edge



years

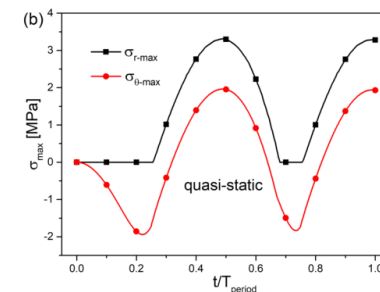
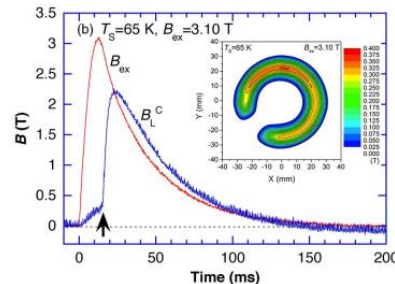
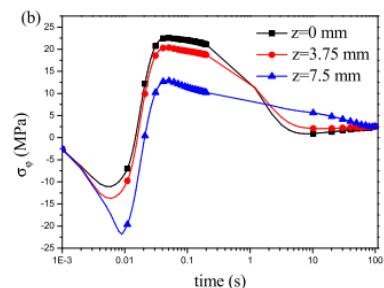
MgB2 bulk

phase-field
damage
during the
flux jump



Stress
distribution

years



Research on the mechanical stability during PFM

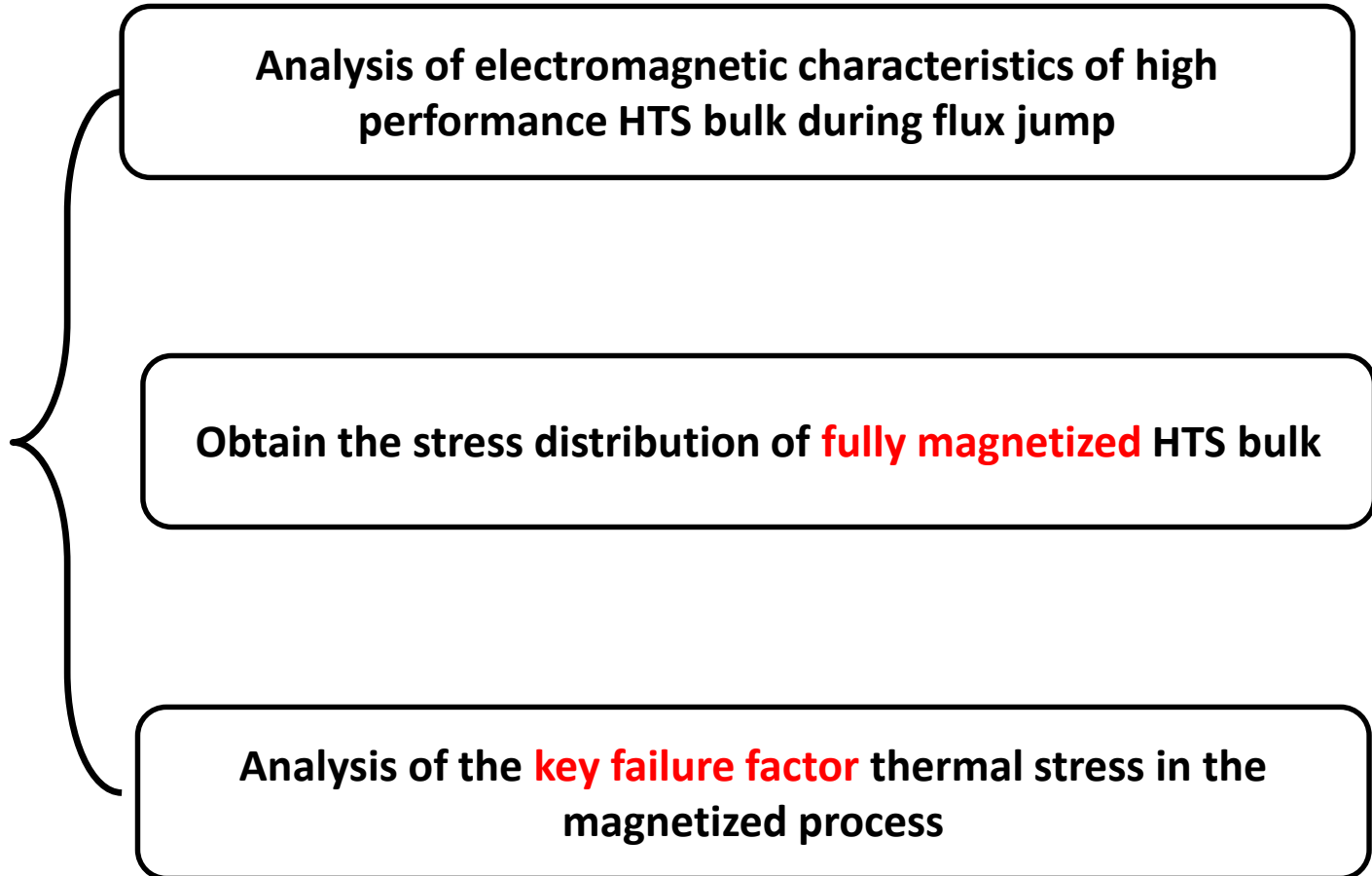
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Author	Sample	Method	Time	respect
Mochizuki	GdBCO bulk	2D Coupled H	2016	coupled thermomagnetic- mechanical instability behavior
Wu	GdBCO bulk	2D Coupled H	2018	
Zhang	MgB2 bulk	2D coupled H	2020	
Hirano	Ring bulks	2D Coupled H	2021	
Shinden	Ring bulks	3D Coupled H	2022	
Ru	GdBCO bulk	2D PD theory	2019	fracture/damage problem
Jing	GdBCO bulk	2D Phase-field theory	2020	
Jing	MgB2 bulk	2D Phase-field theory	2022	

the stronger flux pinning (J_c characteristic) is, the severer the mechanical stability is

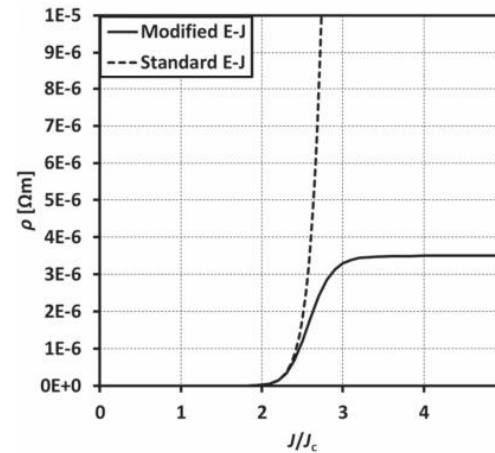
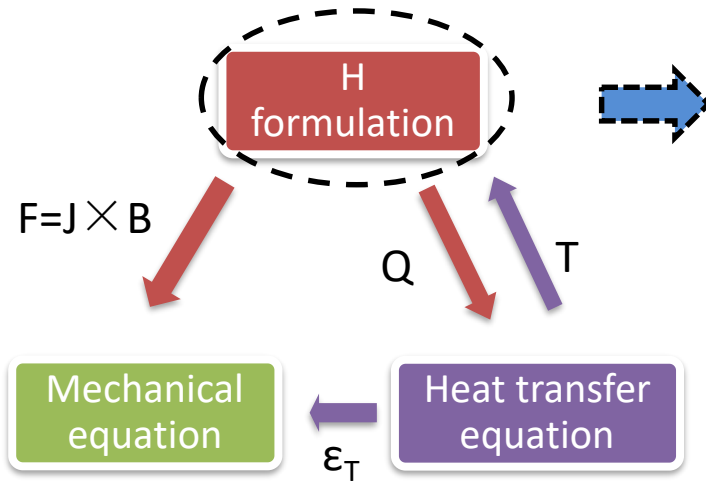


the stronger flux pinning (J_c characteristic) is, the flux jump is prone to occur



- **Numerical model**

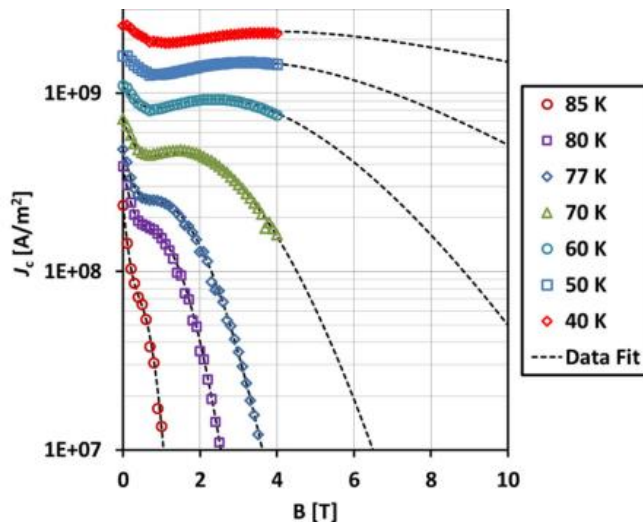
Model setting of magnetic field



Modified E-J relationship

$$\rho_s = \frac{E_c}{J_c(B, T)} \left(\frac{|J|}{J_c(B, T)} \right)^{n-1}$$

$$\rho = \begin{cases} \frac{\rho_s \rho_{\text{norm}}}{\rho_s + \rho_{\text{norm}}}, & T < T_c \\ \rho_{\text{norm}}, & T \geq T_c, \end{cases}$$

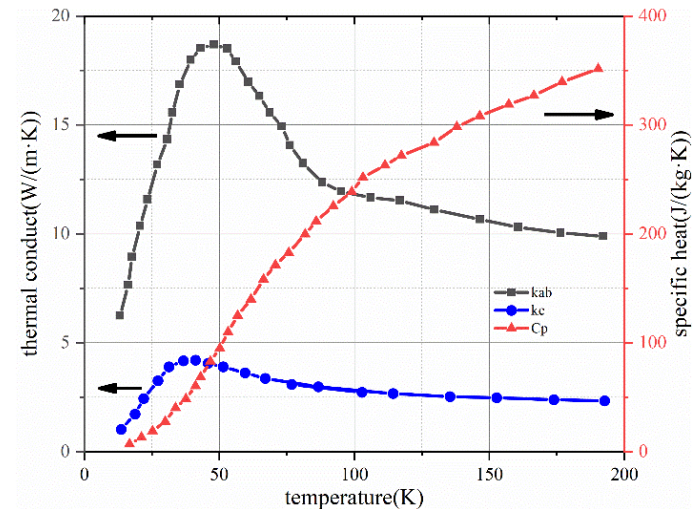
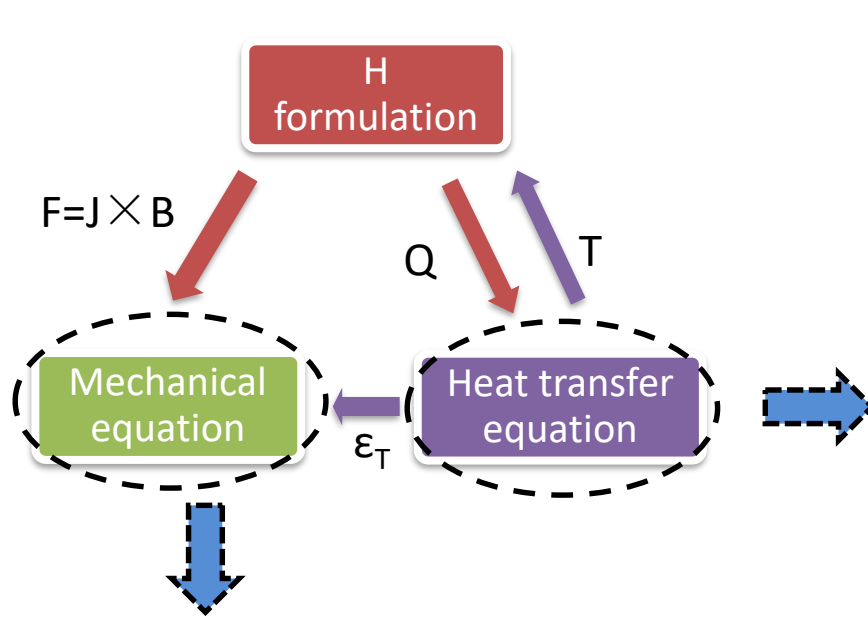


Introduce ρ_{norm} to modify the E-J relationship during flux jump

Adopt high performance HTS bulk J_c expression

Two-variable interpolation J_c characteristic

Model setting of Multiphysics



Key Thermal parameter

$$\nabla \cdot \boldsymbol{\sigma} + \mathbf{f} = 0$$

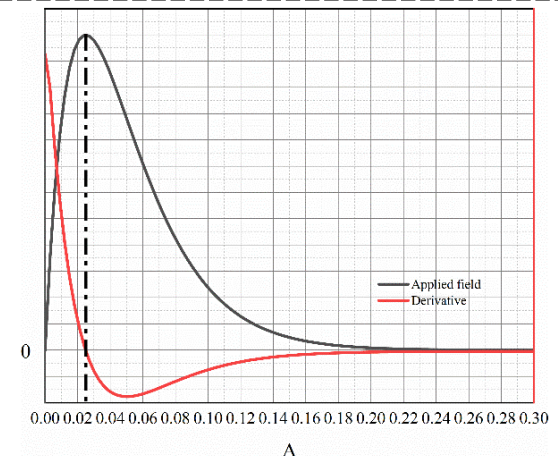
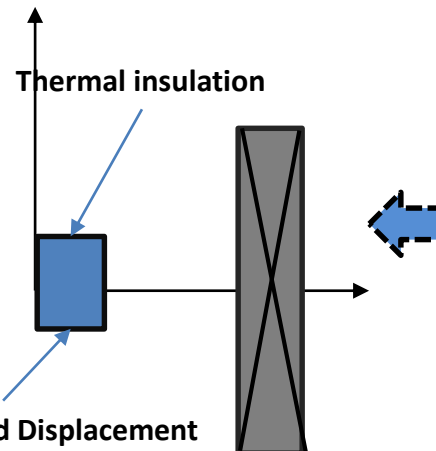
$$\boldsymbol{\epsilon} = \frac{1}{2} [(\nabla \mathbf{u})^T + \nabla \mathbf{u}] + \alpha \cdot \Delta T$$

$$\boldsymbol{\sigma} = \mathbf{D} : \boldsymbol{\epsilon}$$

axial symmetry

Zero Prescribed Displacement

Axial symmetry model



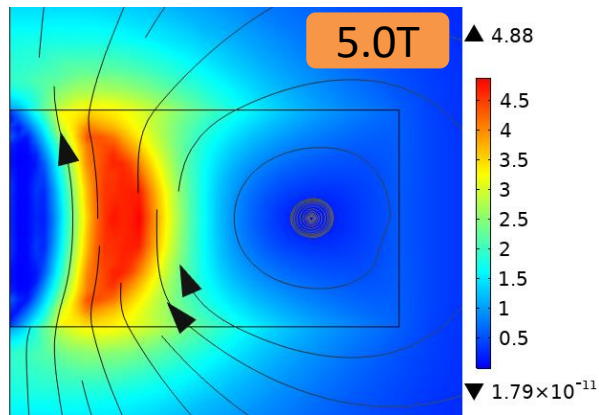
Wave form of applied field



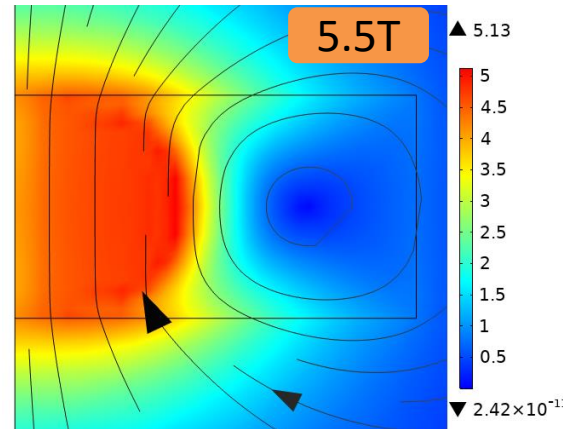
- **Numerical result**

Magnetic distribution after fully magnetized

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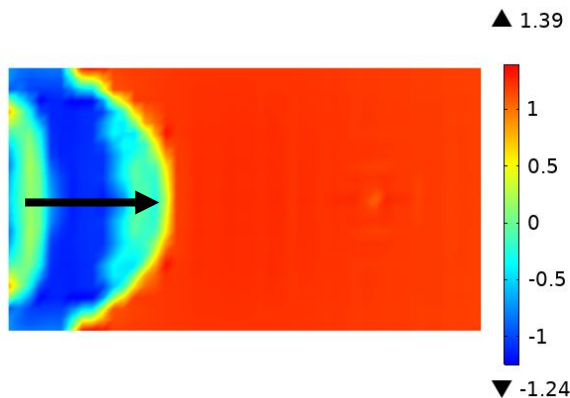


before flux jump

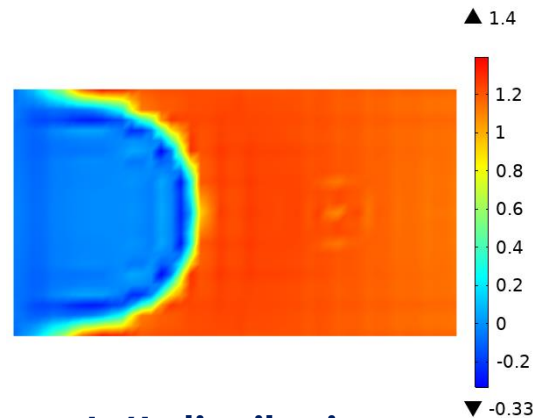


After flux jump

Only increasing 0.5T,
the sample is fully
magnetized



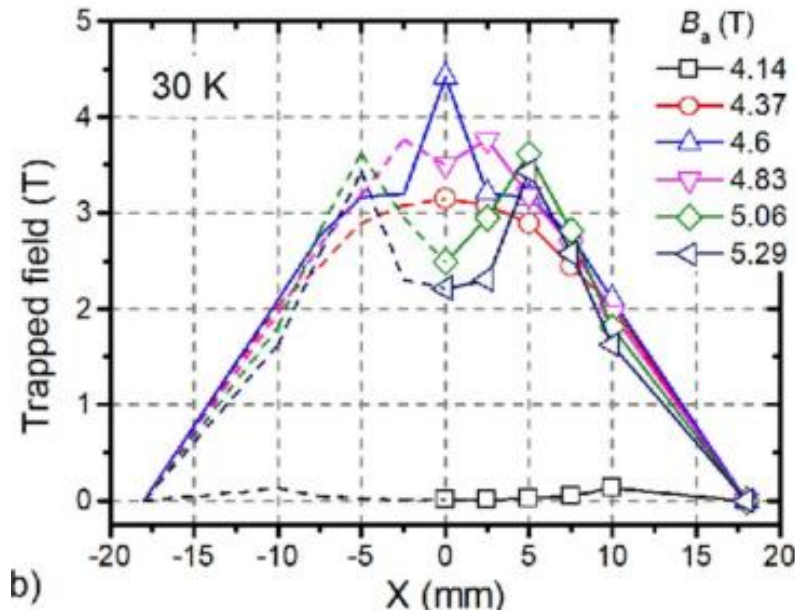
J_ϕ/J_c distribution



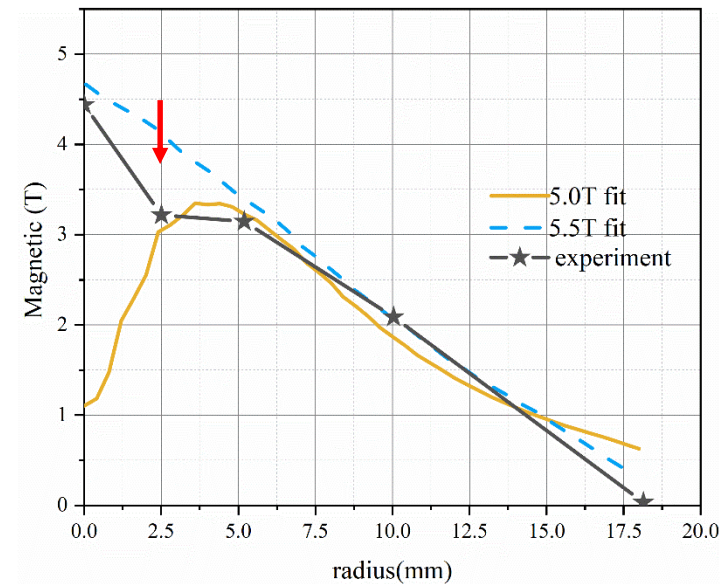
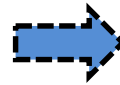
J_ϕ/J_c distribution

After flux jump, the
current is redistributed,
and there is no positive
current at the center

Temperature evolution and distribution



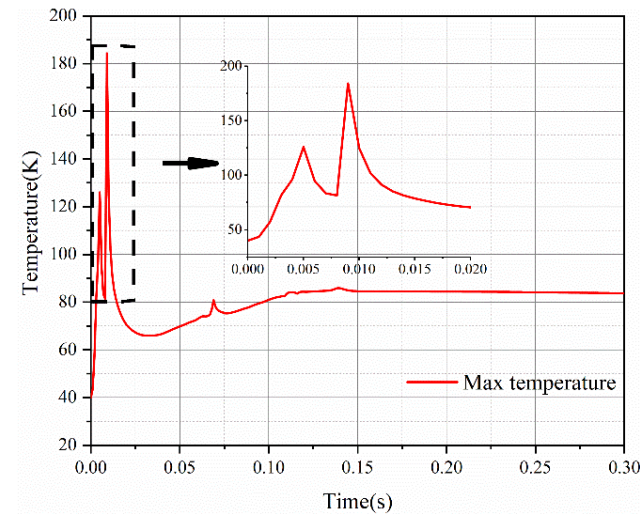
Magnetic distribution of Experiment



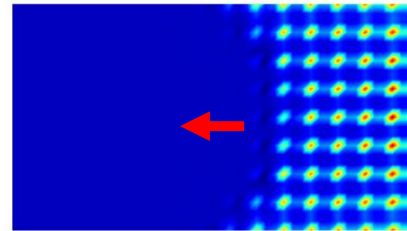
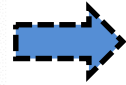
Simulation result

The final trapped field 4.6T has good agreement with the experiment 4.4T, but there is a concave in the position 2.5mm

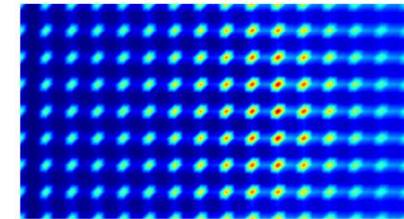
Temperature change during flux jump



Maximum temperature time evolution

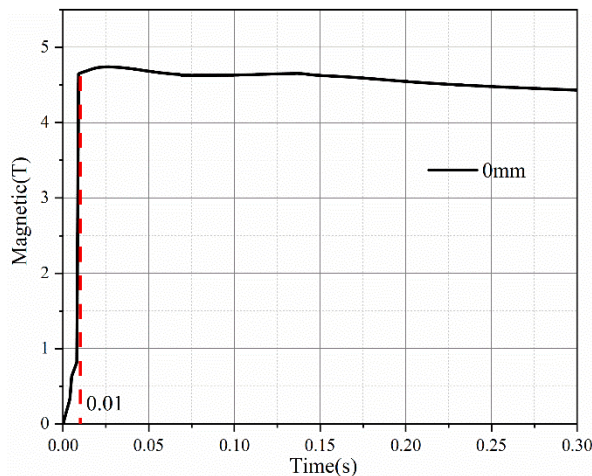


5ms



9ms

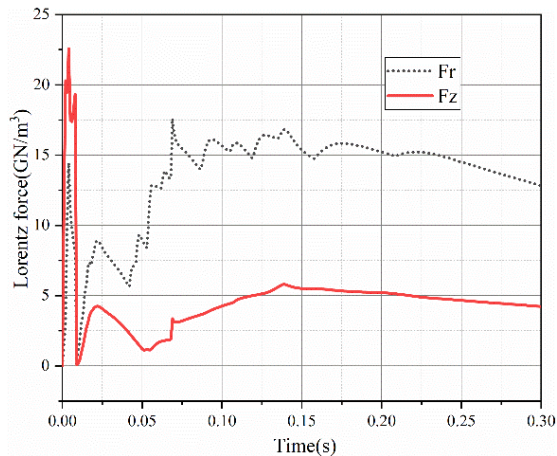
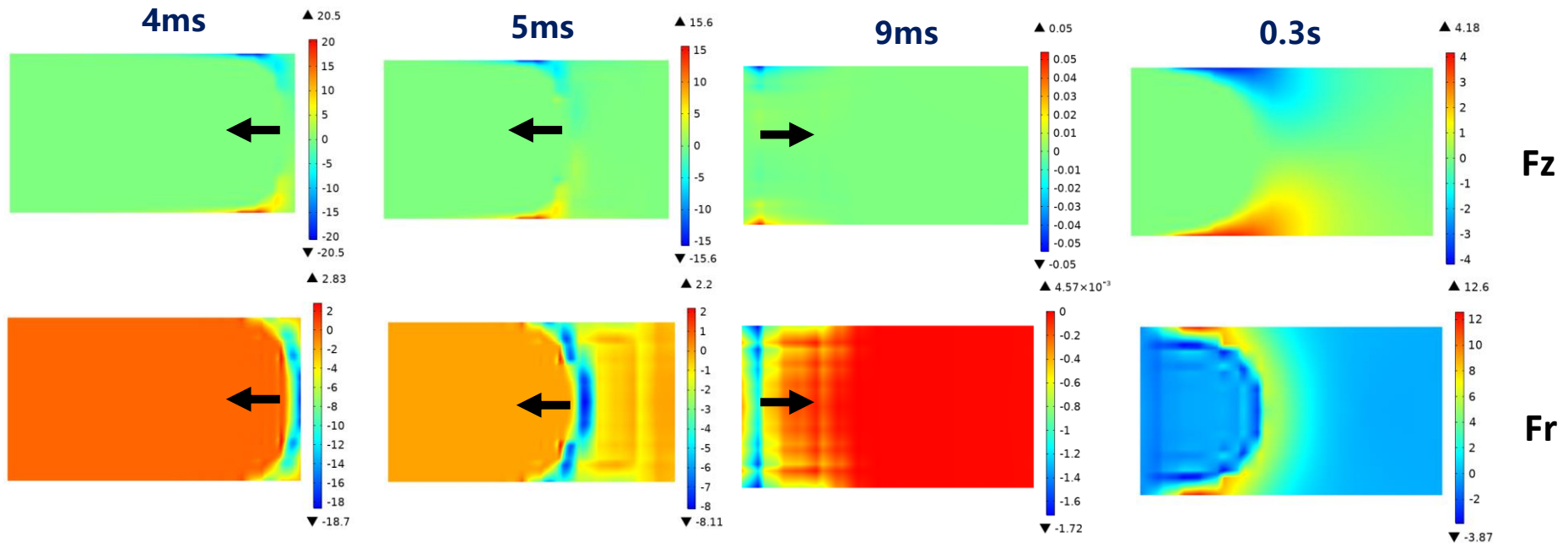
temperature hot pot at peak position



0mm magnetic time evolution

with the temperature attach to the peak, the screen current induced by J_c reduced immediately, resulting in the flux jump

Lorentz force load

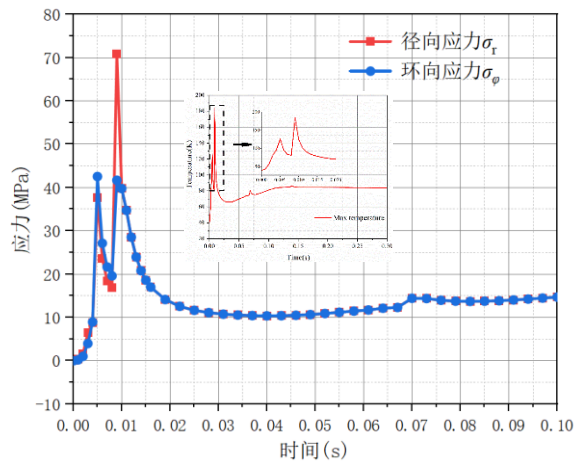
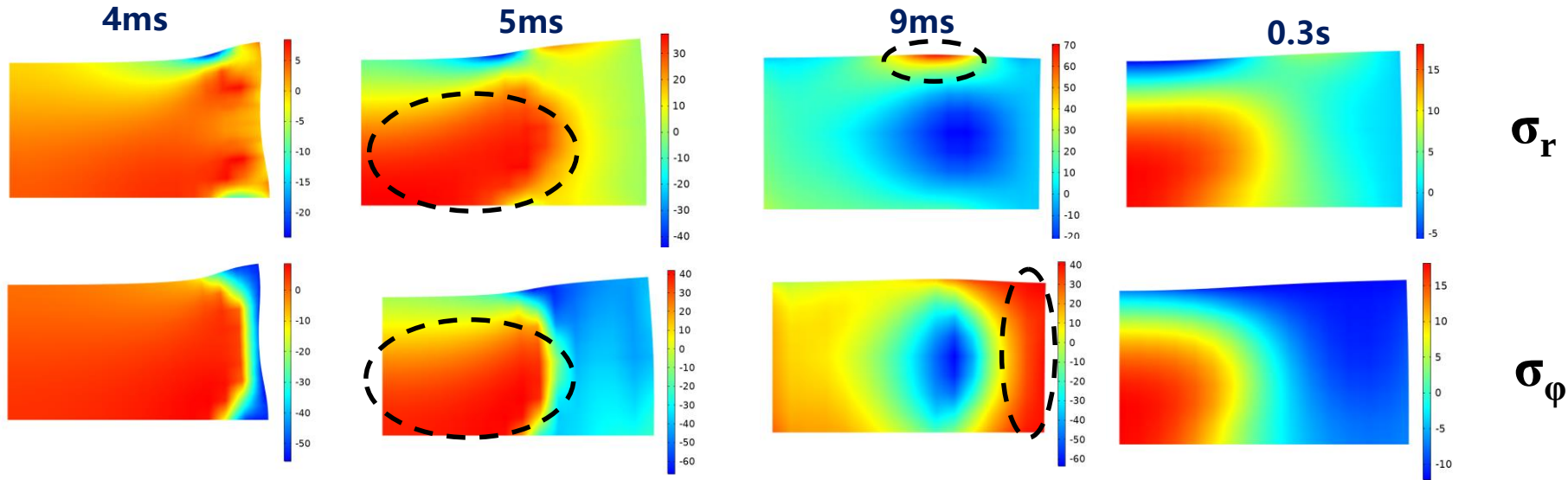


Maximum Lorentz force evolution

the Lorentz forces Move from edge to center and suddenly penetrate into the center, then Slowly return to the edge

Thermal Stress distribution after fully magnetized

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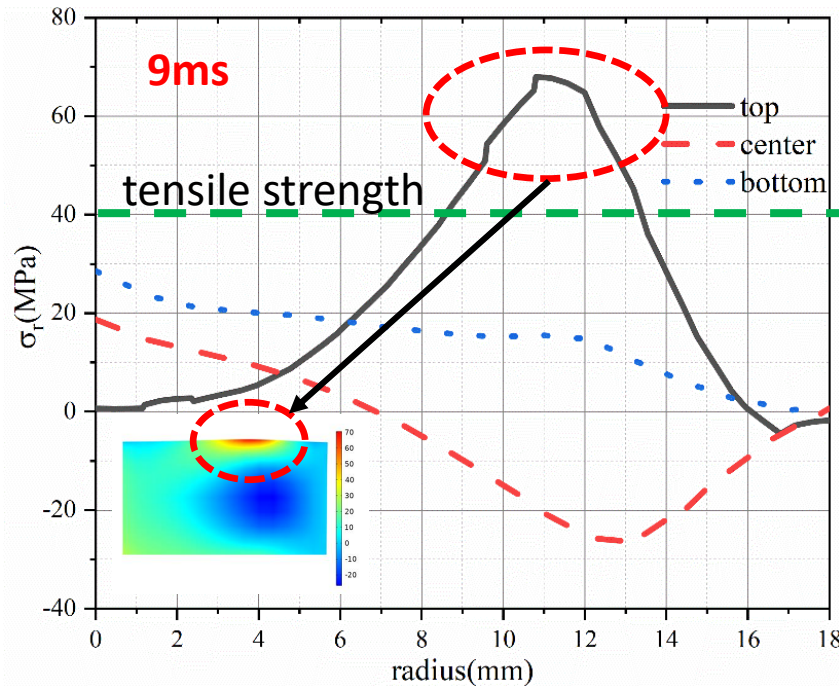


the change of maximum thermal stress is highly consistent with the temperature

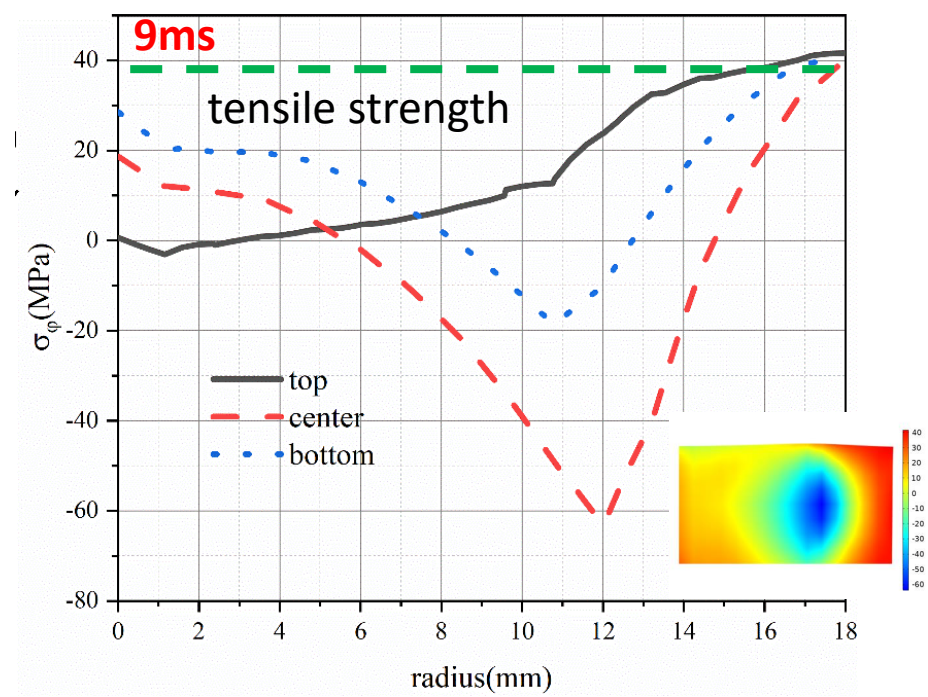
after flux jump, the thermal stress redistributed, and the core stress released, result in local stress concentration

Maximum thermal stress time evolution

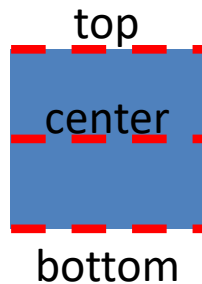
Spatial distribution of thermal stress during flux jump



thermal σ_r Spatial distribution

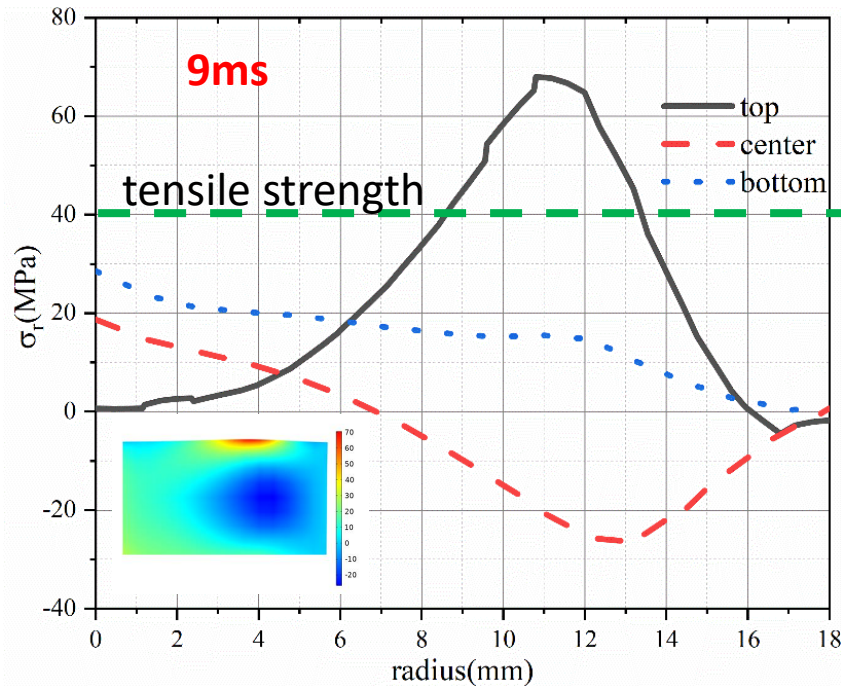


thermal σ_ϕ Spatial distribution

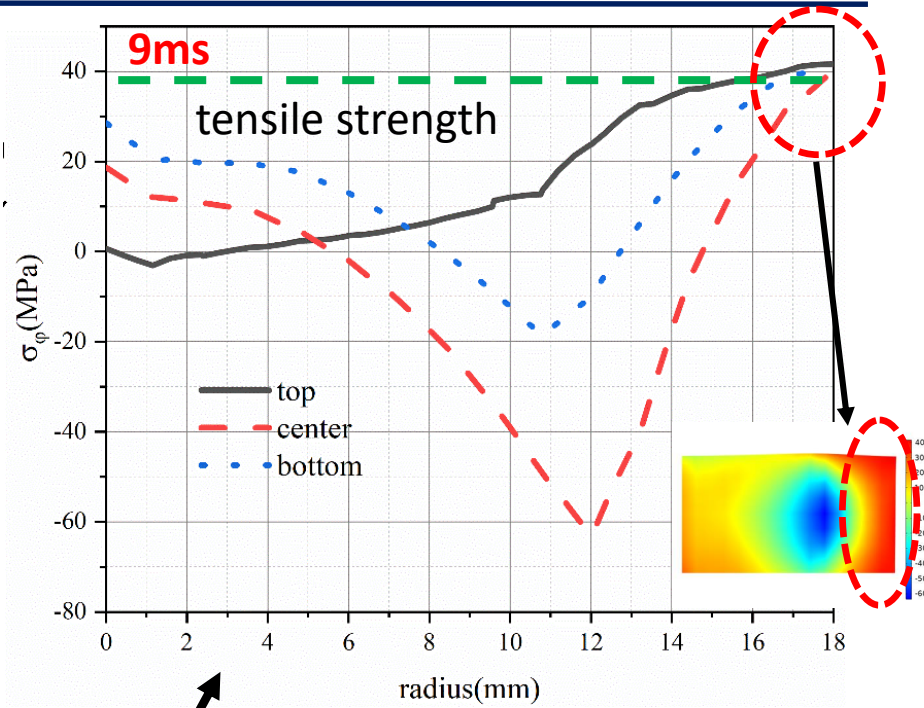


After stress redistribution, the thermal σ_r concentrated at the top surface, then the peak value hits 70MPa and exceeds tensile strength 40MPa, which is dangerous for HTS materials

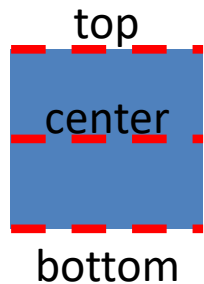
Spatial distribution of stress during flux jump



Thermal σ_r Spatial distribution

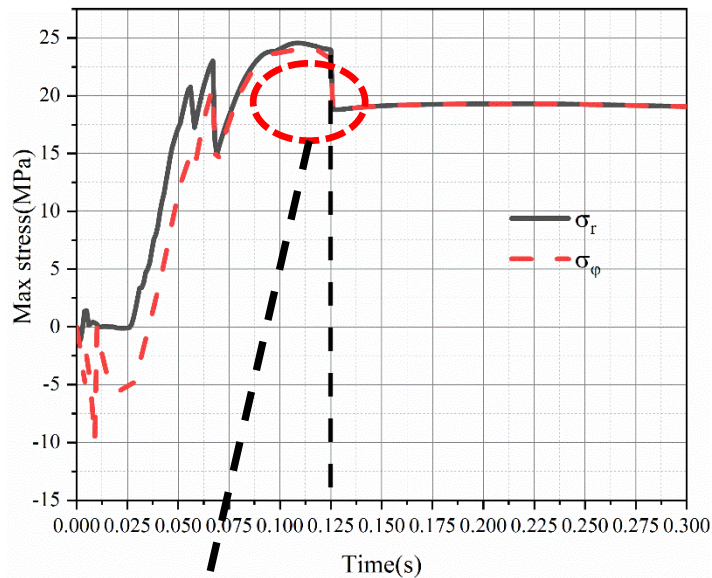


Thermal σ_ϕ Spatial distribution

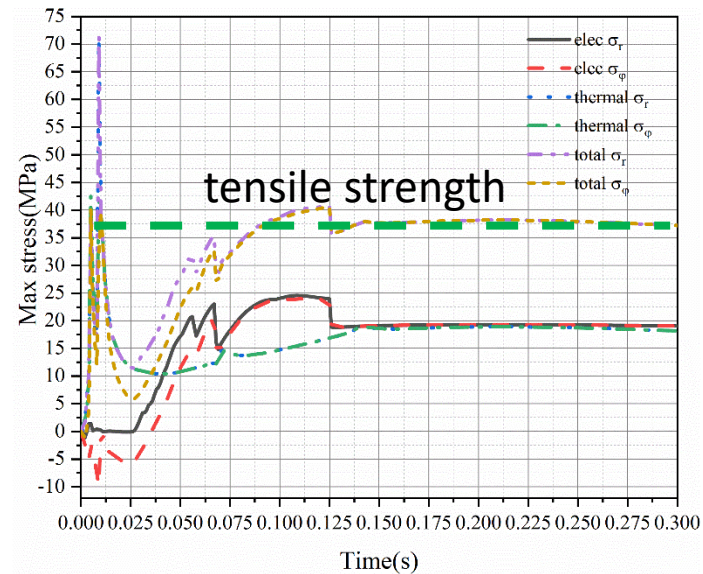


For the σ_ϕ , the stress is concentrated at the outer region of the bulk. the max value is close to the tensile strength 40MPa.

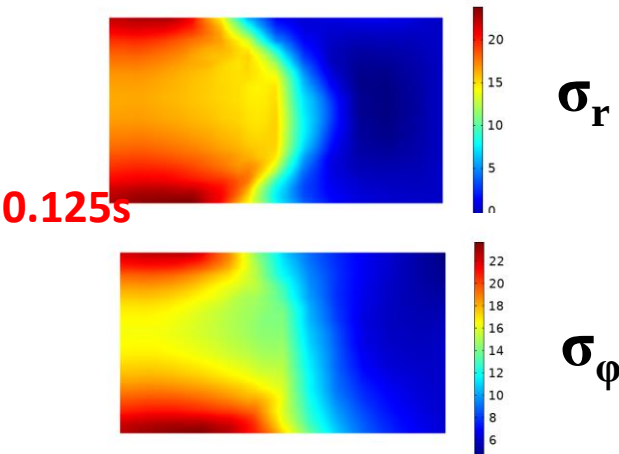
Spatial distribution of electromagnetic stress during flux jump



Max electromagnetic stress time evolution



Max total stress time evolution



Electromagnetic stress is small compared against thermal stress at flux jump, but gradually increases with time goes by

The total stress hit the tensile strength at the end of PFM, which is the obstacle to further improve the trapped field

After flux jump, the electromagnetic and thermal parameters will change rapidly and help the bulk to be fully magnetized

After flux jump, the stress will redistribute to propagate outward, result in local stress concentration

The total stress in fully magnetized bulk trapped 4.8T is close to the mechanical strength, which is the obstacle to further improve the trapped field

THANK YOU